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Charge trapping and retention behaviors of Ge nanocrystals distributed in the gate oxide near the gate synthesized by low-energy ion implantation

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A layer of Ge nanocrystals (nc-Ge) distributed in the gate oxide near the gate of a metal-oxide-semiconductor structure is synthesized with low-energy Ge ion implantation followed by thermal annealing at 800 °C. The behaviors of charge trapping and charge retention in the nc-Ge have been studied. For a positive charging voltage, only electron trapping occurs, and the trapped electrons show a long retention time. However, for a negative charging voltage, both the hole trapping and electron trapping occur simultaneously, and the hole trapping is dominant if the magnitude of the charging voltage is small or the charging time is short. Due to the relatively easier loss of the trapped holes, the net charge trapping in the nc-Ge exhibits a continuous shift toward a more negative value with the waiting time. © 2007 American Institute of Physics.

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I. INTRODUCTION

Si and Ge nanocrystals have a promising application in nonvolatile memory (NVM) devices. 1–5 Although many works have been devoted to Si nanocrystal (nc-Si), Ge nanocrystal (nc-Ge) was believed to have a better performance due to its smaller band gap. 4,5 The nc-Ge embedded in a SiO2 layer can be fabricated with various techniques such as wet oxidation of Si1−xGe x film, 4 chemical vapor deposition, 5,6 cosputtering, 7 and ion implantation. 8,9 In a NVM device based on the nc-Ge, the structure of gate/ control oxide/nc-Ge layer/tunnel oxide/Si substrate is usually formed with the nc-Ge layer located near the Si substrate. 5,10,11 In the present study, the above structure is simplified by using a single low-energy ion implantation into a SiO2 thin film and eliminating the deposition of the control oxide. In this simplified structure, the nc-Ge is distributed in a narrow layer in the gate oxide near the gate rather than the substrate. The simplified structure shows the memory effect and exhibits some interesting behaviors of charge trapping and retention. It is observed that both electrons and holes can be stored in the nc-Ge depending on the polarity and magnitude of the charging voltage as well as the charging time. It is also observed that the electron storage has a much longer retention time as compared to the hole storage.

II. EXPERIMENTAL DETAILS

The fabrication of the metal-oxide-semiconductor (MOS) structure with the nc-Ge distributed near the gate in the gate oxide began with the growth of a 32 nm SiO2 thin film on the p-type Si substrate at 950 °C by dry oxidation. Subsequently, Ge ions were implanted into the SiO2 film with a dosage of 1 × 1016 cm−2 at the energy of 4 keV. According to the stopping and range of ions in matter (SRIM) simulation, a Gaussian distribution (range=7.7 nm and straggle=2.2 nm) of Ge ions in the SiO2 film with the peak (the peak concentration is 1.1 × 1022 at./cm3) located near the SiO2 surface was obtained. The implanted samples were then annealed at 800 °C in N2 ambient for 40 min to induce the formation of nc-Ge. Cross-sectional transmission electron microscopy (TEM) was used to reveal the structural information of the annealed samples. Aluminum was deposited and patterned to form the top and bottom electrodes. At the same time, a control sample without Ge ion implantation was also prepared under identical conditions. To investigate the behaviors of charge trapping and retention in the nc-Ge, the charging experiment and the capacitance-voltage (C-V) measurement were performed in a dark environment at room temperature with a Keithley 4200 semiconductor characterization system and a HP 4284A LCR meter.

III. RESULTS AND DISCUSSIONS

Figure 1 shows the TEM image for the sample with the nc-Ge distributed in the SiO2 thin film. The inset of Fig. 1...
shows the TEM image of a single nc-Ge with a size of about 4 nm. A narrow layer of nc-Ge near the surface of the SiO$_2$ film is clearly observed. The width of the nc-Ge layer is about 10.5 nm, the oxide thickness between the nc-Ge layer and the Si substrate is about 25 nm, and the oxide thickness between the surface of the SiO$_2$ film and the nc-Ge layer is about 4 nm. The total thickness of the SiO$_2$ film containing the nc-Ge is 39.5 nm, which is about 20% larger than the thickness of the unimplanted oxide. This observation consists with the previous studies in which the oxide swellings with thickness of the unimplanted oxide. This observation consists of the trapping of electrons in the nc-Ge distributed in the oxide. Note that the negative sign in Eq. (1) accounts for the fact that the trapping of negative charges in the nc-Ge leads to a positive $\Delta V_{fb}$. As examples, the $Q_{charge}$ estimated with Eq. (1) are $-9.62 \times 10^{-7}$ and $-9.23 \times 10^{-7}$ C/cm$^2$ for $V_G = +25$ and $-25$ V, respectively.

To study the influence of both the polarity and magnitude of the charging voltage on the charge trapping in the nc-Ge, the $Q_{charge}$ was measured for both positive and negative $V_G$ with a magnitude varying from 7 to 30 V at the fixed charging time of 1 s. The results are shown in Fig. 3. It is evident from Fig. 3 that the charging behavior for positive $V_G$ is different from that for negative $V_G$. As shown in the figure, a positive charging voltage always results in a negative $Q_{charge}$ regardless of the voltage magnitude. In contrast, a negative charging voltage leads to either a positive or a negative $Q_{charge}$ depending on the magnitude of the voltage. In the case of positive $V_G$, for $V_G$ less than $+21$ V, the $Q_{charge}$ remains at a small negative value with a magnitude less than $1 \times 10^{-7}$ C/cm$^2$, indicating that only a small amount of electrons are trapped in the nc-Ge. However, when the $V_G$ goes beyond $+21$ V, the $Q_{charge}$ is still negative, but its magnitude increases rapidly (for example, $Q_{charge}$ = $-27 \times 10^{-7}$ C/cm$^2$ for $V_G$ = +30 V) and is approximately a linear function of the $V_G$. As will be discussed later, the rapid increase in the magnitude of $Q_{charge}$ is due to the strong electron injection from the Si substrate under the influence of a large positive gate voltage. On the other hand, in the case of negative $V_G$, the $Q_{charge}$ is positive when the magnitude of $V_G$ is less than 19 V, and the maximum $Q_{charge}$ ($3.2 \times 10^{-7}$ C/cm$^2$) occurs at $V_G$ = -16 V. When the magnitude of $V_G$ is larger than 19 V, the $Q_{charge}$ becomes negative and its magnitude rapidly increases as the magnitude of the $V_G$ is further increased. It is obvious that in the case of negative $V_G$ both hole trapping
and electron trapping in the nc-Ge occur. As will be discussed later, under a negative \( V_G \), the hole trapping and electron trapping are two competing processes, and the magnitude of the \( V_G \) determines which one is dominant. This behavior is different from the situation of a conventional memory structure where the nc-Ge layer is located near the Si substrate.\(^9\) In a conventional structure, the electron trapping caused by a negative \( V_G \) is insignificant because the thick control oxide prevents the electron injection from the gate.

The results shown in Fig. 3 indicate that the positive \( V_G \) results in only electron trapping, whereas the negative \( V_G \) leads to both electron trapping and hole trapping simultaneously. As shown in the inset of Fig. 3(a), under the influence of a positive \( V_G \), electrons are injected from the Si substrate and trapped in the nc-Ge. When the magnitude of \( V_G \) is larger than 21 V, the electric field in the oxide between the Si substrate and the nc-Ge layer is higher than 5 MV/cm, leading to the occurrence of Fowler-Nordheim (FN) tunneling across the oxide. Therefore, the electron injection is drastically enhanced, resulting in a rapid increase in the electron trapping in the nc-Ge. On the other hand, as shown in the inset of Fig. 3(b), under the influence of a negative \( V_G \), the injection of the holes from the Si substrate and the injection of electrons from the gate occur at the same time, leading to the trapping of both holes and electrons in the nc-Ge. In this case, the density of the net trapped charges in the nc-Ge is given by

\[
Q_{\text{charge}} = Q_{\text{hole}} - Q_{\text{electron}},
\]

where the \( Q_{\text{hole}} \) and \( Q_{\text{electron}} \) are the trapped-hole density and trapped-electron density, respectively. The behavior shown in Fig. 3(b) is a consequence of the competition between these two charging processes. Obviously, \( Q_{\text{charge}} \) is positive if the hole trapping dominates, and it is negative if the electron trapping dominates. Figure 3(b) indicates that in the case of a negative \( V_G \), when its magnitude is less than 18 V, the hole trapping is dominant, but the electron trapping increases rapidly when the magnitude of \( V_G \) is larger than 16 V. The rapid increase of the electron trapping is due to the occurrence of electron tunneling from the gate electrode. When the magnitude of the negative \( V_G \) is larger than 18 V, the electron tunneling from the gate surpasses the hole injection from the substrate, leading to a net negative charge trapping in the nc-Ge. This situation is in contrast to that of a conventional memory device where the nc-Ge is located near the substrate. For the conventional memory device, a negative \( V_G \) is usually used during the erasing process to expel the trapped electrons from the nc-Ge, and the negative \( V_G \) itself does not cause the electron trapping in the nc-Ge because of the thick control oxide which can effectively suppress the electron tunneling from the gate (note that only the trapping of positive charges in the nc-Ge is observable in the case of over-erasing).

Further investigation on the charging mechanisms is carried out by varying the charging time from 5 \( \mu \)s to 5 s for a given \( V_G \), and the result is shown in Fig. 4. The result consists with the above-mentioned charging mechanisms. As shown in Fig. 4(a), the \( Q_{\text{charge}} \) is always negative regardless of the charging time for a positive \( V_G \). This is because only electron injection from the substrate occurs under a positive \( V_G \). It can be observed in Fig. 4(a) that for the \( V_G = +25 \) V, the electron trapping in the nc-Ge starts to increase rapidly when the charging time is longer than 1 ms. This starting time is prolonged when the \( V_G \) is reduced. Obviously, this is related to the \( V_G \) dependence of the electron injection from the substrate. On the other hand, as shown in Fig. 4(b), for the charging by a negative \( V_G \), both positive and negative \( Q_{\text{charge}} \) could be observed depending on the magnitude of the \( V_G \) and the charging time. As discussed earlier, under a negative \( V_G \), the two competing processes (i.e., the hole injection from the substrate and the electron injection from the gate) occur simultaneously. For \( V_G = -15 \) and -18 V, the \( Q_{\text{charge}} \) is positive (i.e., the hole trapping is dominant) and increases gently with the charging time. The increase of the \( Q_{\text{charge}} \) with the charging time for \( V_G = -15 \) V is even faster than that for \( V_G = -18 \) V. This is because the electron injection under \( V_G = -15 \) V is less significant. For \( V_G = -20 \) V, a small amount of positive charge trapping could be observed for a charging time shorter than 1 ms, but the electron trapping becomes dominant when the charging time is longer than 1 ms. For \( V_G = -23 \) and -25 V, the electron trapping is dominant, and the amount of the \( Q_{\text{charge}} \) increases with the charging time. When the charging time is longer than 1 s, the
electron trapping is saturated. Such phenomenon is due to the strong Coulomb interaction between the trapped electrons in the nc-Ge and the negatively biased gate electrode (note that the nc-Ge layer is close to the gate electrode).

The difference in charging process between positive and negative $V_G$ has an impact on the charge retention. The charge retention characteristics of the memory structure are shown in Fig. 5. The charge retention experiment was carried out for both positive $V_G$ [Fig. 5(a)] and negative $V_G$ [Fig. 5(b)] with a fixed charging time of 1 s. For the charging under a positive $V_G$, as discussed above, only electrons are trapped in the nc-Ge, and thus the $Q_{\text{charge}}$ behavior in the charge retention experiment is simple, i.e., the $Q_{\text{charge}}$ remains negative and a decrease in its magnitude could be observed. As can be seen in Fig. 5(a), for both the charging voltages of +18 and +25 V, the $Q_{\text{charge}}$ remains unchanged for the waiting time up to $10^3$ s, and then only a slight loss of the trapped electrons is observed. The result indicates that the two oxide layers sandwiching the nc-Ge layer, as shown in the inset of Fig. 5(a), have a good capability to prevent the release of the trapped electrons. However, the charge retention behavior for a negative charging voltage is quite different. As shown in Fig. 5(b), for the $V_G$=−18 V, the $Q_{\text{charge}}$ is positive when the waiting time is shorter than $10^3$ s, but later it becomes negative and shifts continuously toward a more negative value (i.e., its magnitude increases) with the waiting time. For the $V_G$=−25 V, the $Q_{\text{charge}}$ is always negative regardless of the waiting time, and it also shifts continuously toward a more negative value with the waiting time. However, the rate of the change in $Q_{\text{charge}}$ for $V_G$=−25 V is lower than that for $V_G$=−18 V. As discussed earlier, both electrons and holes are trapped in the nc-Ge under a negative charging voltage. However, the above results indicate that the loss of the trapped holes is much faster than that of the trapped electrons. Therefore, with the faster loss of the trapped holes, the $Q_{\text{charge}}$ can change from positive to negative as observed in the case of charging by $V_G$=−18 V, and it shifts continuously toward a more negative value with time. As the hole trapping for $V_G$=−18 V is more significant than that for $V_G$=−25 V, the rate of the change in $Q_{\text{charge}}$ for $V_G$=−18 V is larger than that for $V_G$=−25 V as mentioned above. The faster loss of the trapped holes could be due to the fact that the trapped holes can easily recombine with the electrons moving from the Al gate electrode.

IV. CONCLUSION

In summary, we have synthesized a narrow layer of nc-Ge located near the gate in the gate oxide of the MOS structure using the low-energy Ge ion implantation and the
thermal annealing at 800 °C. The charge trapping and charge retention in the nc-Ge have been studied. For a positive charging voltage \( V_G \), only electron trapping occurs, which is due to the electron injection from the \( p \)-type substrate. The electron trapping depends strongly on the magnitude of the \( V_G \) and the charging time. It increases rapidly with the \( V_G \) when the FN tunneling across the oxide between the nc-Ge layer and the substrate occurs. However, for a negative \( V_G \), both the hole trapping due to the hole injection from the substrate and the electron trapping due to the electron injection from the gate occur, and these two competing mechanisms determine whether the net charge trapping in the nc-Ge is positive or negative. The hole trapping could be dominant if the magnitude of \( V_G \) is small and the charging time is short, and the electron trapping becomes dominant when the magnitude of \( V_G \) is large and the charging time is sufficiently long. For a positive \( V_G \), the electron trapping shows good retention. In contrast, for a negative \( V_G \), due to the easier loss of the trapped holes, the net charge trapping in the nc-Ge shows a continuous shift toward a more negative value with the waiting time, and the shift strongly depends on the magnitude of the \( V_G \) and the charging time.

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