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Influence of pulsed laser deposition rate on the microstructure and thermoelectric properties of Ca$_3$Co$_4$O$_9$ thin films

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Abstract

The effects of deposition rate on the microstructure and thermoelectric properties of Ca$_3$Co$_4$O$_9$ thin films fabricated by pulsed laser deposition (PLD) technique were investigated. X-ray diffraction (XRD) and high-resolution transmission electron microscopy (HRTEM) revealed that a fast deposition rate resulted in not only low crystallinity but also the existence of the Ca$_x$CoO$_2$ secondary phase. Formation of Ca$_x$CoO$_2$ was inevitable during the thin film growth, and this was discussed from both structural and compositional point of view. With longer deposition interval or with sufficient oxygen at a lower deposition rate, the Ca$_x$CoO$_2$ phase was able to transit into the desired Ca$_3$Co$_4$O$_9$ phase during the coalescence process. The quality of the thin films was further analyzed by electrical properties measurements. The Ca$_3$Co$_4$O$_9$ thin film fabricated at a
slower deposition rate was found to exhibit a low electrical resistivity of 9.4 mΩcm and high Seebeck coefficient of 240 μV/K at about 700°C, indicating a good quality film.

**Keywords:**

**PACS Classification numbers:**
61.05.cp X-ray diffraction
68.37.Ps Atomic force microscopy (AFM)
72.80.Ga Transition-metal compounds
73.50.Lw Thermoelectric effects

1. Introduction

Thermoelectric (TE) materials can directly convert heat energy into electricity power or vice versa. Therefore, they are of great interests for application in cooling and power generation devices [1]. The performance of a TE material is usually characterized by the dimensionless figure of merit $ZT = S^2 \sigma T / \kappa$, where $S$, $\sigma$, $T$ and $\kappa$ are the Seebeck coefficient, electrical conductivity, absolute temperature, and thermal conductivity, respectively.

Recently, misfit layered cobalt oxides have been identified to exhibit good thermoelectric properties and have received considerable attention thereafter due to their advantages over intermetallic alloys in terms of high temperature stability and low toxicity [2, 3]. Ca$_3$Co$_4$O$_9$ (CCO) is among the best misfit layered cobalt oxides. Due to its good thermal stability at 1000K, intensive studies have been made to date. CCO crystal is composed of stacks of CdI$_2$-type CoO$_2$ layer alternating with rock-salt-type Ca$_2$CoO$_3$ layer along the c-axis and can be denoted as $[\text{Ca}_2\text{CoO}_3]^{1.62}[\text{CoO}_2]_{1.62}$ [3]. The $ZT$ value obtained for a single crystal reached ~0.83 at 1000 K [4]. However, it is very difficult to grow single crystals with large sizes. In addition, randomly
oriented ceramics show a relatively poor TE performance because CCO exhibits a high anisotropy in its TE properties.

Thin films, especially epitaxial ones, offer strongly oriented structure in contrast to bulk ceramics, and can exhibit intrinsic electrical properties similar to bulk single crystals. Moreover, thin film structure also benefits $ZT$ improvement by increasing the thermopower due to quantum confinement effects and reducing the thermal conductivity due to lattice mismatches and scattering at surfaces and interfaces [5, 6]. CCO epitaxial thin films have been fabricated by RF sputtering and pulsed laser deposition (PLD) techniques on MgO (100), sapphire (0001) and Si (100) substrates [7-9]. However, thin films with high crystal quality have not been prepared and high temperature TE properties were also seldom reported. Moreover, an understanding on the CCO thin film growth mechanisms and the existence of secondary phase in the thin film structure [8, 9] are also limited. Although a number of PLD processing parameters have been investigated, the deposition rate, which is one of the most important parameters during thin film deposition process, has not received proper attention in earlier works on CCO thin film preparations. Here, we report our investigation of the deposition rate on CCO thin films, which indicated that a relatively low deposition rate is necessary to achieve a well-crystallized CCO structure as well as good thermoelectric properties. A faster deposition rate, on the other hand, results in the presence of the $\text{Ca}_x\text{CoO}_2$ secondary phase in the thin films. The growth mechanism of the CCO thin films will be discussed and explanation on the formation and further reaction of the secondary phase will also be presented from both the structural and chemical point of view.
2. Experimental Procedures

A $\text{Ca}_3\text{Co}_4\text{O}_9$ ceramic plate, synthesized by conventional solid state reaction, was used as the target for CCO thin films growth. Prior to the deposition, sapphire (0001) (c-plane) substrates were cleaned in ethanol and water followed by annealing at 1400°C in oxygen gas flow to achieve an atomically-flat surface. The deposition was carried out by KrF excimer laser beam (Lambda Physik Compex, $\lambda$=248 nm) in 20 Pa oxygen pressure and the substrates were heated to 750°C. The deposition rate could be related to many parameters such as laser energy density and target distance. In this work, the laser energy intensity was fixed at 1.2 J/cm$^2$ per pulse and the distance between substrate and target was 5 cm. The deposition rate was adjusted by varying the laser repetition frequency since this could be independently varied without giving rise to any side effects. The laser repetition frequency of 3, 4, 5 and 8 Hz were chosen, which correspond to deposition rates of ~2 to 5 nm/min.

The crystal quality, orientation, thickness and the surface morphology of the thin films were evaluated by X-ray diffraction (XRD), surface profiler, cross-sectional high-resolution transmission (HRTEM), and atomic force microscopy (AFM). The thin films thickness were about 100, 108, 80 and 110 nm for samples deposited using 3, 4, 5, and 8 Hz laser repetition respectively. XRD measurements were conducted using X’Pert PRO (PANalytical) with monochromatic Cu Kα radiation ($\lambda$=1.5406 Å) within 5-75° 2θ range. HRTEM observations were performed using a JEOL 2100F at 200 kV equipped with an Energy Dispersive X-ray (EDX) detector (Oxford Instruments, UK). For 3-Hz and 5-Hz-deposited thin films, TEM specimens were fabricated using a Focused Ion Beam dual beam system (Helios Nanolab 600 or Carl Zeiss...
Nvision 40) and finally cleaned by using fine beam milling at 5 and 2 kV. Surface morphologies of CCO thin films were analyzed using a MFP-3D Stand Alone AFM (Asylum Research) at room temperature. The in-plane electrical resistivity ($\rho$) of the thin films was measured by dc four-probe method and the in-plane Seebeck coefficient ($S$) was measured by a conventional steady state method from room temperature to ~700°C in helium atmosphere.

3. Results and discussion

Figure 1 shows the high-resolution XRD results of the CCO thin film deposited using 3 Hz laser repetition, and the inset shows the $\omega$-scan rocking curve recorded around the (002) diffraction peak. The strong $c$-axis texture for the thin film revealed that the CCO thin film was heteroepitaxially grown on the (0001) plane of $\alpha$-Al$_2$O$_3$ substrate over the entire area of the film. From the seven {001} peaks of the film, the $c$-axis parameter was calculated to be 10.74(6) Å, which is in reasonable agreement with the bulk value (10.83 Å) [3] and other reported thin films (10.7 Å) [8-10]. The small full width at half maximum (FWHM) of the rocking curve (0.06°) is much smaller than reported value (0.83°) [11], which indicates that the films are well crystallized, i.e. of high crystallinity. Note that, by using the 3 Hz laser repetition, the deposition rate was only about 2 nm/min. This is a fairly low value as compared to normal PLD of other oxides which varies from 2 to 15 nm/min [12, 13]. If the thin films were prepared at a higher deposition rate using higher laser repetition, the intensities of their XRD peaks decrease significantly as shown in Figure 2. In particular, for the thin films prepared using 5 Hz, peaks corresponding to the (001) and (003) planes are very weak while peaks corresponding to the (002) and (004) planes maintain high intensities with a slight shift to lower $2\theta$ values. This XRD pattern is found to be similar to
that of the layer-structured Ca$_x$CoO$_2$ [14], which are made up of alternating Ca cation and CoO$_2$ layers.

Surface morphology of the as-prepared thin films was further analyzed using an AFM scanning a 5×5 μm square area (Figure 3). All three AFM images show tiny islands on the surface of the samples prepared using 3, 4 and 5 Hz laser repetition. It is obvious that for 3-Hz-deposited thin film, the island size is smaller and a smoother surface with a root mean square roughness (RMS) of ~4.2 nm was achieved. Meanwhile, by using faster deposition rates larger islands appeared on the sample surfaces, and a higher roughness (RMS: 6~8 nm) was obtained. These island-surface morphologies indicate that the thin film growth mechanism could be ascribed to either the Volmer-Weber growth mode (or island-growth mode) or the Stranski-Krastanov growth mode (S-K growth mode), which are common for hetero-epitaxial thin film growth [15]. For the S-K growth mode, the thin film growth mechanism starts with a layer-by-layer mode in the first stage, and this is followed by an island-growth mode after the critical thickness. The island formation on the surface of our thin films indicates a relatively high thermodynamic wetting energy of CCO film material on the sapphire substrate which could be ascribed to the large lattice mismatch (1.5% for CoO$_2$ layer/Al$_2$O$_3$ and 4.5% for the Ca$_x$CoO$_3$ layer/Al$_2$O$_3$) [11]. Since the nuclei growth or islands coalescence is a kinetic process in the latter stage of the thin film growth, a low deposition rate always caters for a good thin film quality with high crystallinity and smooth surface, while fast deposition rate results in high density of defects, random orientation of nuclei or even amorphous structure. It is obvious that our XRD and AFM results agree well with the above trends.

To further investigate the crystal structure of the thin films, especially the structure of the secondary phase, HRTEM analysis were performed with the incident beam parallel to the [1100]
of sapphire substrate on the cross-sectional area. Figure 4 (a) and (b) shows a well-ordered pile-up of layered structures in the 3-Hz-deposited CCO thin film. The d-spacing was evaluated to be 10.7 Å, which is consistent with the XRD analysis. The epitaxial relation of the film along the c-axis on the c-plane sapphire was found to be [001]Ca$_3$Co$_4$O$_9$//[0001]Al$_2$O$_3$. A transition region about 5 nm thick existed between the crystalline CCO thin film and the sapphire substrate. This region was found to consist of both secondary phase and amorphous structures. The secondary phase was adjacent to the surface of the substrate and was identified to be the Ca$_x$CoO$_2$ structure with a c-plane distance nearly half that of CCO. A small amount of stacking faults was also observed in the CCO structure. Nevertheless, the CCO structure still maintained good crystallinity on both sides of the stacking fault as shown in Figure 4 (c).

Figure 5 represents the HRTEM images of the 5-Hz-deposited thin film. Figure 5(b) shows another layered structure with 5.4 Å lattice d-spacing, which corresponds to the Ca$_x$CoO$_2$ phase [14]. The (001) plane of the Ca$_x$CoO$_2$ phase is still parallel to the (0001) plane of the sapphire substrate. However, such clear layered structure was only found in the region about 10 nm away from the film-substrate interface. Further away from the interface, a large amount of stacking faults were observed in the 5-Hz-deposited thin film, and the microstructure consists of a mixture of CCO, Ca$_x$CoO$_2$ and amorphous structure as shown in Figures 5(c) - (e). The images clearly show that the structure cannot maintain as either CCO or Ca$_x$CoO$_2$ on both sides of the stacking faults, but instead result in the formation of an amorphous phase (Figure 5c), curved layered lattice (Figure 5d) or even structures with other layered orientation (Figure 5e). All these correspond to a bad quality thin film. Ca$_x$CoO$_2$ or amorphous structure among CCO structures have also been reported in other CCO thin film fabricated by sputtering or PLD technique [8, 9]. However, no explanations were provided.
No layer-by-layer CCO growth region could be found at the first few atom layers near the substrate interface of the 3-Hz and 5-Hz deposited samples observed under HRTEM. Moreover, for the 3-Hz deposited films grown at a low deposition rate, CCO structures grew only after the transition layer, which was of mixed composition and a few atom layers thick. Herein, the CCO thin film growth could be further confirmed as an island-growth mechanism. The formation of the transition layer could be ascribed to the high wetting energy at the substrate-film interface and the high super-saturation during the PLD process. The stress from the lattice mismatch was totally released within the transition layer. The texture of CCO thin film was then achieved only by a self-assembling process without any epitaxial relationship from the substrate. Such observation was in good agreement with recent reports that highly textured CCO thin films were successfully fabricated on both Si (001) and glass substrate [9, 16]. Therefore, the quality of the CCO thin film is more related to the coalescence process where nuclei or islands unite while the atoms access to their thermodynamically stable sites.

From the structure point of view, similar to the structure of Na$_x$CoO$_2$ (0.3≤ x ≤1), the crystal structure of Ca$_x$CoO$_2$ is built up by alternating Ca cation and CoO$_2$ layers. Its structure is compared with the CCO structure in Figures 6(a) and (b). When the thin films stack up, it is easier to get a Ca$_x$CoO$_2$ structure since the c-axis lattice parameter of Ca$_x$CoO$_2$ is only about 5.4 Å [14], which is nearly half that of CCO (10.8Å). Meanwhile, the CoO$_2$ layer, composed of edge-shared CoO$_6$ octahedra, is a fundamental sub-unit in both CCO and Ca$_x$CoO$_2$ crystal structures. Due to the strong Co-O bonding, CoO$_2$ layer may be formed preferentially as compared to the Rock-salt layer. Thus, at stacking fault sites CoO$_2$ layers tend to extend further and a Ca$_x$CoO$_2$ structure is favored to release the local defect energy as illustrated in Figure 6(c). On the other hand, from the composition point of view, Ca$_x$CoO$_2$ may have a wide range of Ca:Co elemental
ratio, but contains lower O:Co ratio of 2:1 as compared to that of CCO which is 2.25:1. Thus, another possible reason for the Ca₄CoO₂ formation may lie in the lack of oxygen locally. Given enough time and sufficient oxygen supply, part of the CoO₂ layers are inclined to combine dissociative Ca and O atoms into the rock-salt layer of Ca₂CoO₃ and the Ca₄CoO₂ structure is able to transform into the CCO structure for lower thermodynamic energy as shown in Figure 6(d). The transition from Ca₄CoO₂ into CCO structure is a feasible process and has been reported elsewhere for the preparation of CCO [10, 17]. Nevertheless, a possible CCO thin film growth mechanism can be proposed as follows: During the coalescence process of nuclei, it is inevitable to have Ca₄CoO₂ structure at nuclei boundaries with high density of defects, or at some regions lacking in oxygen. Given enough ripening time, the defects can move to the surface of the thin film when top-layer nuclei unit together. The Ca₄CoO₂ can also combine with the ambient oxygen to transform into CCO structure. Both processes would result in a well aligned CCO thin film with c-axis texture. However, if the deposition rate is too high, too much mass will reach the surface within a unit time and the density of defects will be higher. Thus, more stacking faults will be embedded and accumulated at the boundaries, which will result in the generation of more Ca₄CoO₂ regions. Moreover, the opportunity to combine with ambient oxygen is quite small since the cations are buried under the new arriving layers and isolated from the ambient oxygen atmosphere. So, these embedded Ca₄CoO₂ structure are not able to transit into CCO and will remain within the thin film eventually. In addition, post annealing was also performed on the 5 and 8 Hz-deposited thin films at 800ºC in oxygen flow for 2 hours. This did not improve the CCO structure but instead increased the crystallinity of the Ca₄CoO₂ phase in the thin films. Thus, it is further confirmed that the embedded Ca₄CoO₂ structure can hardly transit into CCO and this
is consistent with reported fact that once formed, $\text{Ca}_x\text{CoO}_2$ is stable up to temperatures higher than 800ºC [17].

Finally, the electrical properties of the thin films deposited at 3, 4 and 5 Hz were compared and shown in Figure 7. Figure 7(a) shows the temperature dependence of the resistivity $\rho$ for the thin films. Within the temperature range investigated, the 3-Hz-deposited thin film exhibits the lowest electrical resistivity, which ranges from 4.6 to 9.4 m$\Omega$cm, while the 4-Hz-deposited sample shows a slightly higher value. The room temperature resistivity of around 5 m$\Omega$cm obtained for these two thin films is fairly low among the values reported for both CCO thin films and single crystals [4, 10, 16], which indicates the good quality of these two thin films. Similar Seebeck coefficients $S$ were also obtained for these two thin film samples. As shown in Figure 7(b), $S$ value increased from 110 $\mu$V/K at room temperature to 240 $\mu$V/K at ~700ºC, which are comparable to reported data (110~240 $\mu$V/K) [4, 10, 16]. The slight difference in the electrical properties ($\rho$ and $S$) between the two samples could be ascribed to the larger amount of stacking faults or lower crystallinity in the thin film prepared at a faster deposition rate. The thin film prepared by using 5 Hz repetition exhibited an electrical resistivity of 15~40 m$\Omega$cm and Seebeck coefficient of 120~180 $\mu$V/K Seebeck coefficient over the temperature range investigated. These properties correspond to a very bad quality thin film. According to its microstructure, the existence of the $\text{Ca}_x\text{CoO}_2$ structure could significantly affect the electrical properties of the thin film, since $\text{Ca}_x\text{CoO}_2$ exhibits a higher $\rho$ and lower $S$ as compared to CCO [18]. Also, defects and amorphous structures could lead to high scattering effects and reduction of carrier concentration, which both result in the higher resistivity obtained in the 5-Hz deposited thin films.
The presence of the structural defect was said to be helpful in enhancing the overall TE performance by reducing the thermal conductivity [16]. And it is indeed feasible to control the density of stacking faults in the CCO thin film by simply using different deposition rate. However, high density of stacking faults can also result in the formation of Ca$_3$CoO$_2$ phase with poor TE properties. Hence, additional study is still necessary to achieve the optimum amount of such stacking faults to obtain the best TE properties in CCO thin films using PLD.

4. Conclusion

In summary, high quality thin films of misfit layered cobalt oxide Ca$_3$Co$_4$O$_9$ were successfully fabricated by pulsed laser deposition with a deposition rate of ~2 nm/min using 3 Hz laser repetition. The XRD and HRTEM results revealed that a fast deposition rate leads to not only low crystallinity but also the formation Ca$_3$CoO$_2$ structures in the thin film. It was also concluded that the quality of the Ca$_3$Co$_4$O$_9$ thin film is strongly related to the coalescence process at the latter stage of the thin films growth. Slow deposition can restrain the formation of Ca$_3$CoO$_2$ secondary phase or help to convert the Ca$_3$CoO$_2$ into the CCO structure. Ca$_3$Co$_4$O$_9$ thin films fabricated by 3 and 4 Hz laser repetition showed electrical resistivity varying from 4.6 to 9.4 m$\Omega$cm and Seebeck coefficient varying from 110 to 240 $\mu$V/K over the investigated temperature ranging from room temperature to ~700ºC. These values were comparable to reported values and indicated that the films were of good quality.
Acknowledgments

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References


Figure Captions

Figure 1. High-resolution XRD patterns of Ca$_3$Co$_4$O$_9$ thin film deposited by using 3 Hz laser repetition on c-plane sapphire substrate. The $\omega$-scan rocking curve around the (002) diffraction peak is shown in the inset.

Figure 2. XRD patterns comparison of the thin films deposited using different laser repetition within a 5º~55º 2$\theta$ range using a high intensity XRD scanning.

Figure 3. Surface morphologies of the thin films prepared at (a) 3, (b) 4 and (c) 5 Hz laser repetition obtained using AFM scanning over a 5x5 µm area.

Figure 4. HR-TEM cross-section images of the Ca$_3$Co$_4$O$_9$ thin film deposited using 3 Hz laser repetition. The inset of (a) is the electron-diffraction pattern of a selected area with good crystallinity. Although stacking faults were observed in the thin film (c), the structures on both sides of the stacking faults still exhibit good crystallinity. (Note: Pt is a protective layer deposited to prevent damaging the CCO thin film during the TEM specimen preparation process)

Figure 5. HR-TEM cross-section images of the Ca$_3$Co$_4$O$_9$ thin film deposited using 5 Hz laser repetition. The inset of (b) is the fast Fourier transform image of a selected area of the secondary phase. Great amount of stacking faults resulted in (c) low crystallinity or amorphous structure, (d) layer curvature and (e) random orientation. (Note: SiO$_2$ is a protective layer deposited to prevent damaging the CCO thin film during the TEM specimen preparation process)

Figure 6. Crystal structures of (a) Ca$_3$Co$_4$O$_9$ and (b) Ca$_3$CoO$_2$, and the illustrations of (c) Ca$_x$CoO$_2$ formation instead of Ca$_3$Co$_4$O$_9$ to reduce the stacking fault energy, and (d) structure transition from Ca$_3$CoO$_2$ to Ca$_3$Co$_4$O$_9$ with combination of ambient oxygen during thin film growth.
Figure 7. Temperature dependence of electrical resistivity (a) and Seebeck coefficient (b) of the thin films deposited using 3, 4, and 5 Hz laser repetition.
Figure 1
Figure 2
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Figure 6
Figure 7